

# **FUNDAMENTALS OF ELECTRON MICROSCOPY THEORY**

**NRAM PRACTICAL COURSE**

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# THE ELECTRON MICROSCOPE HAS RECOGNIZABLE OPTICAL PARTS

- **ELECTRON “GUN”**  
[equivalent to a light source]
- **CONDENSOR LENS SYSTEM**
- **SPECIMEN STAGE**
- **OBJECTIVE LENS**
- **“PROJECTOR LENSES”**
  - FURTHER MAGNIFY THE IMAGE,
  - OR RELAY AN IMAGE OF THE DIFFRACTION PATTERN THAT IS PRODUCED IN THE FOCAL PLANE OF THE OBJECTIVE LENS

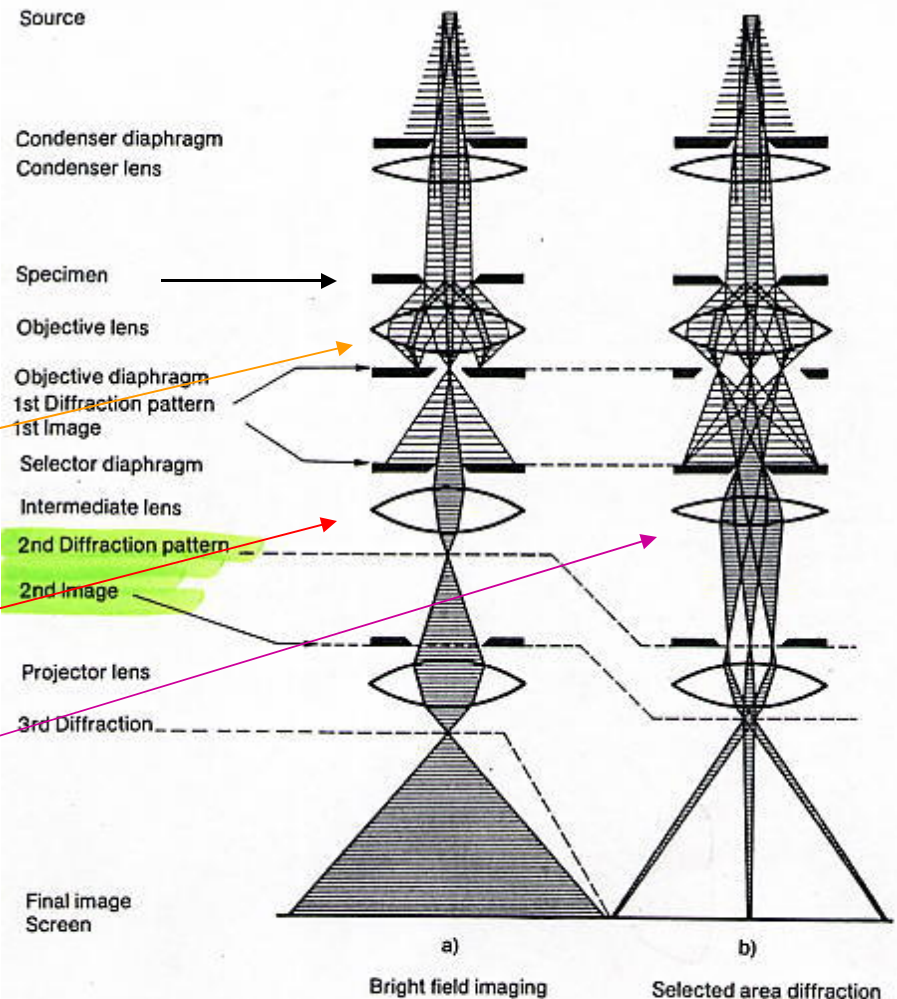
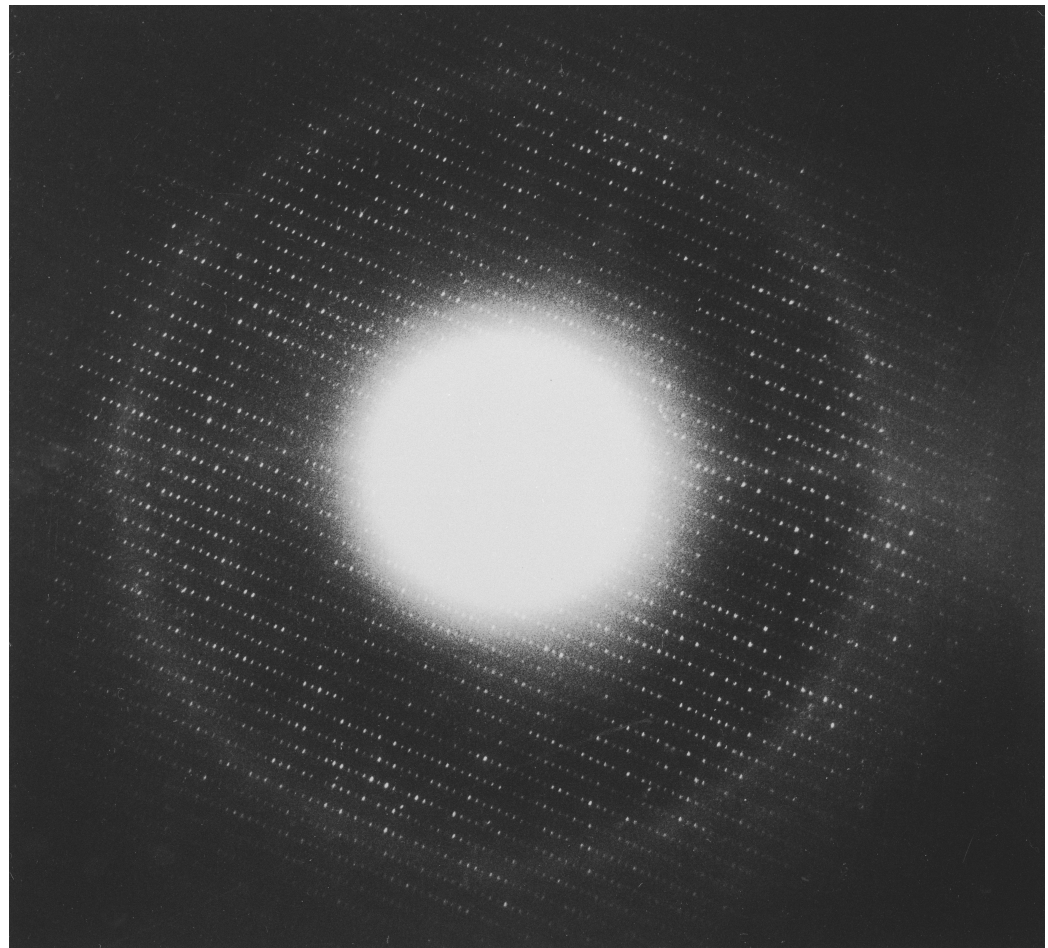


Fig. 4.20 a, b. Ray diagram for a transmission electron microscope in (a) the bright-field mode and (b) selected-area electron diffraction (SAED) mode

# ELECTRONS REALLY ARE WAVES

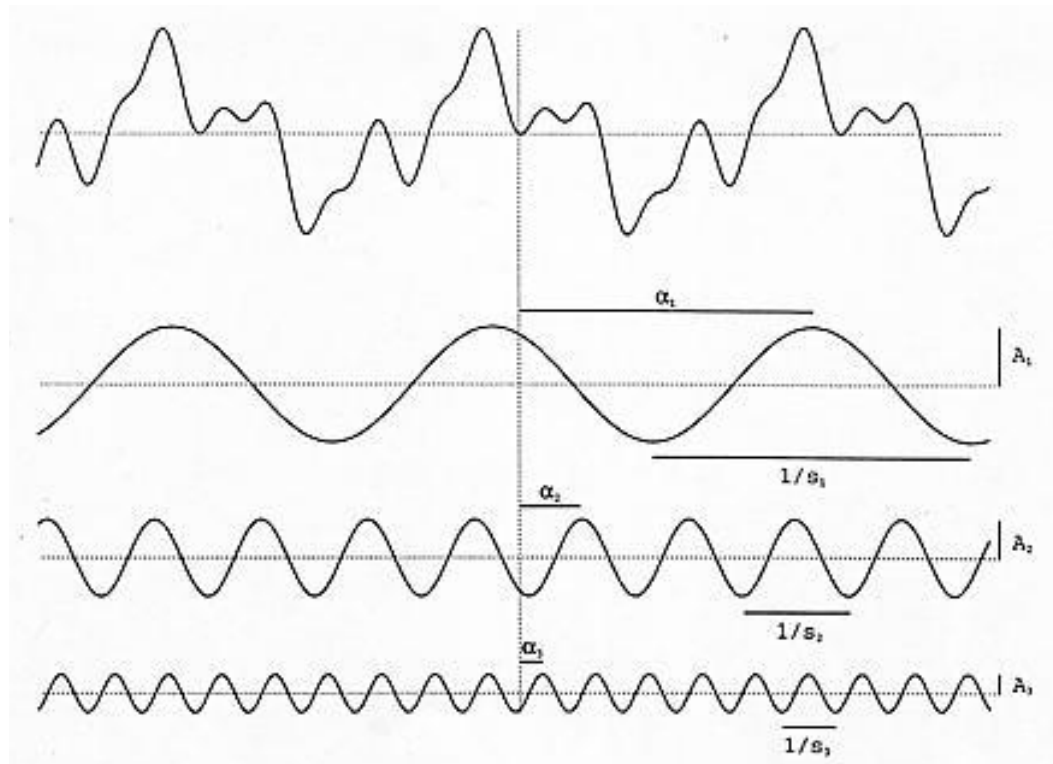
## — AND DIFFRACTION IS IMPORTANT IN EM

- Electrons produce diffraction patterns
  - just like those produced by x-rays
- Lens aberrations and defocus produce phase contrast
  - even though the intensity transmitted through the specimen is *almost* constant
- **Heads up - electrons are also a flux of ionizing radiation ...**



Electron Diffraction Pattern of Catalase

# EACH SCATTERED BEAM IN THE DIFFRACTION PATTERN CONTRIBUTES A SINE-FUNCTION IN THE IMAGE

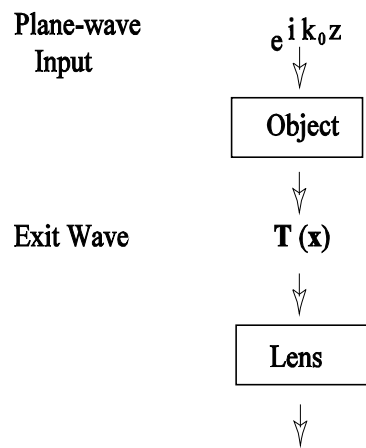
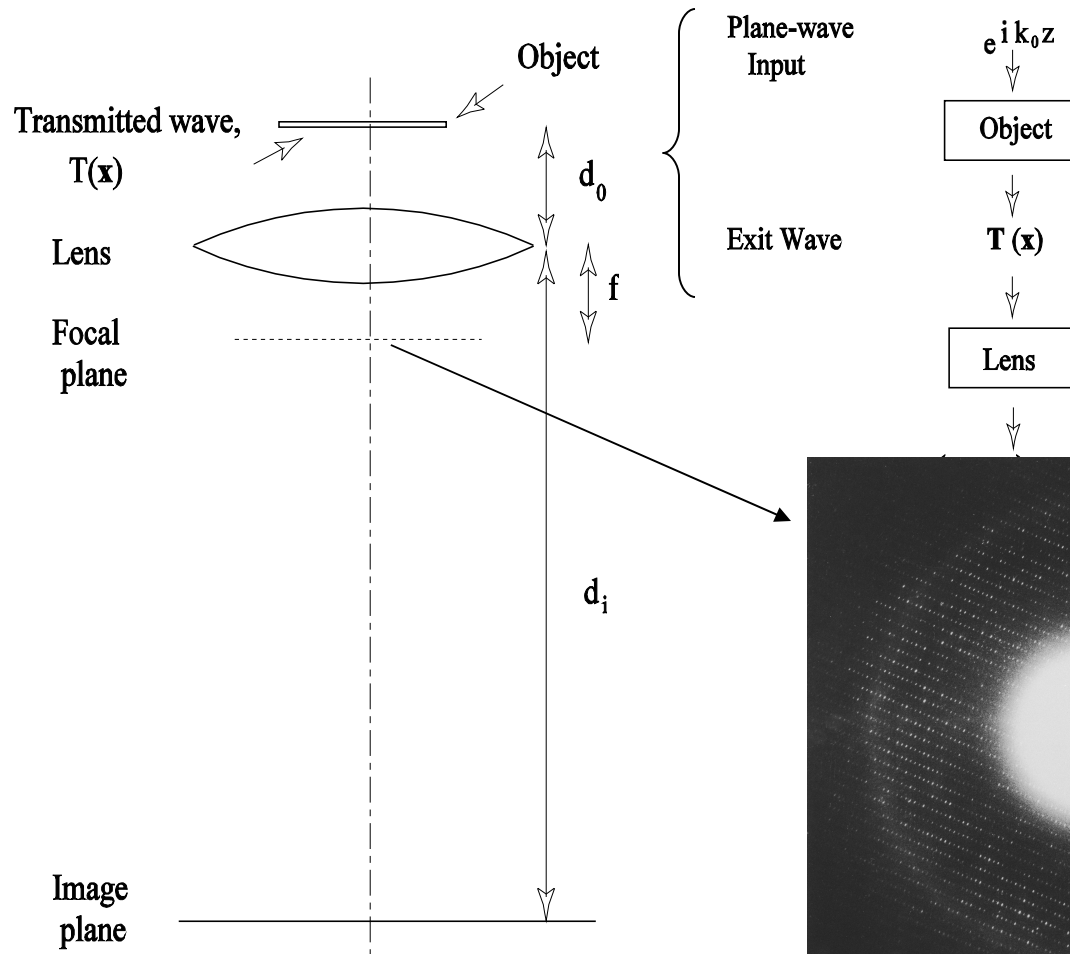


Chiu et al. (1993)  
Biophys J. 64:1610-1625

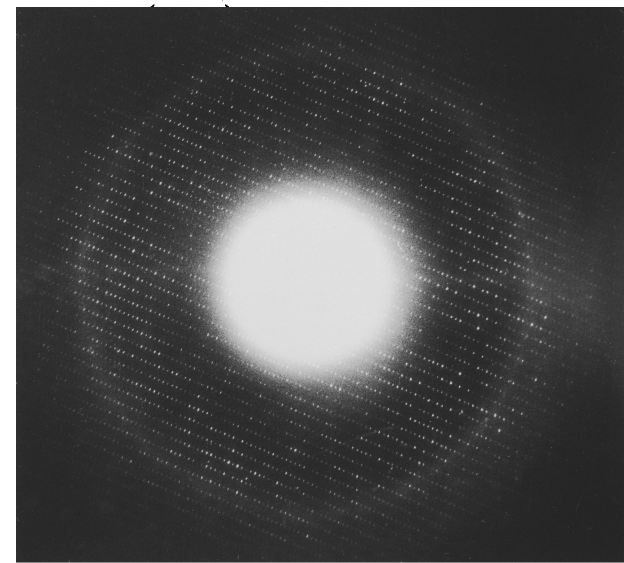
- Each sine-function has its own amplitude and phase
  - Larger scattering angles correspond to higher resolution
- The sine-functions add up to give a complicated function
  - e.g. the image of a molecule
- Crystals help to explain these concepts
  - but everything remains the same when there is no crystal



# THE SCATTERED ELECTRON WAVE FUNCTION IS THE FOURIER TRANSFORM OF THE TRANSMITTED ELECTRON WAVE



The Fourier transform, i.e.  $F(T(\mathbf{x}))$ , is simply a “list of the values of the amplitudes and the phases for every sine function that makes up the transmitted wave”

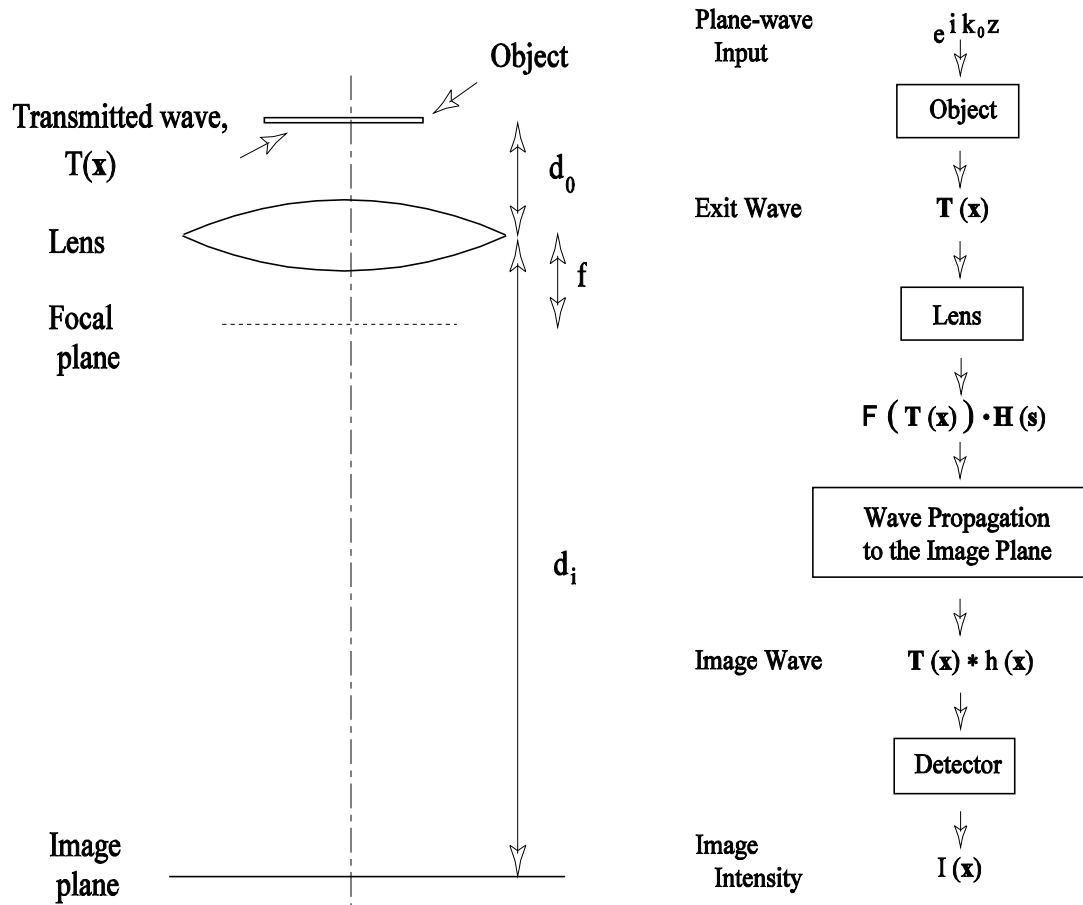


# ABBE'S THEORY OF IMAGE FORMATION

## APPLIES TO ELECTRON WAVES JUST AS IT DOES TO LIGHT

- The scattered wave is the **Fourier transform** of the wave function transmitted through the object
- The lens of a microscope inevitably applies some **aberration function,  $H(s)$** , to the scattered wave
- The wave function in the image is the ***INVERSE* operation (inverse Fourier transform)**
  - But now the inverse step is applied to the *aberrated* wave function, so the result is not the same as the original, transmitted wave
- The **image intensity is the square** of the image wave function

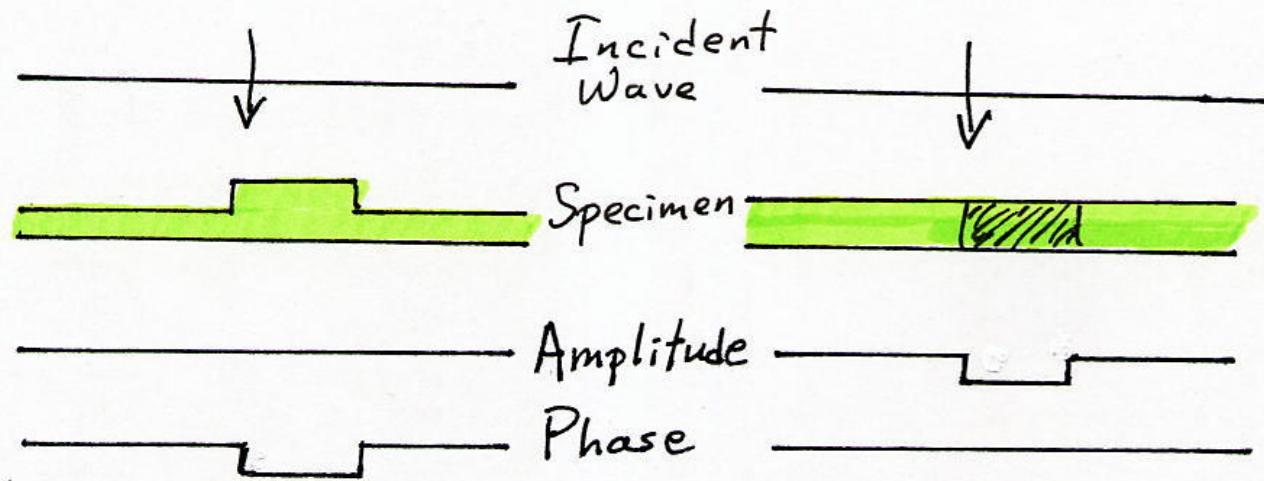
# THE IMAGE WAVE IS THE INVERSE FOURIER TRANSFORM OF THE SCATTERED (AND ABERRATED) ELECTRON WAVE



$H(\mathbf{s})$  represents the wave aberration (and the effect of a limited lens-aperture)

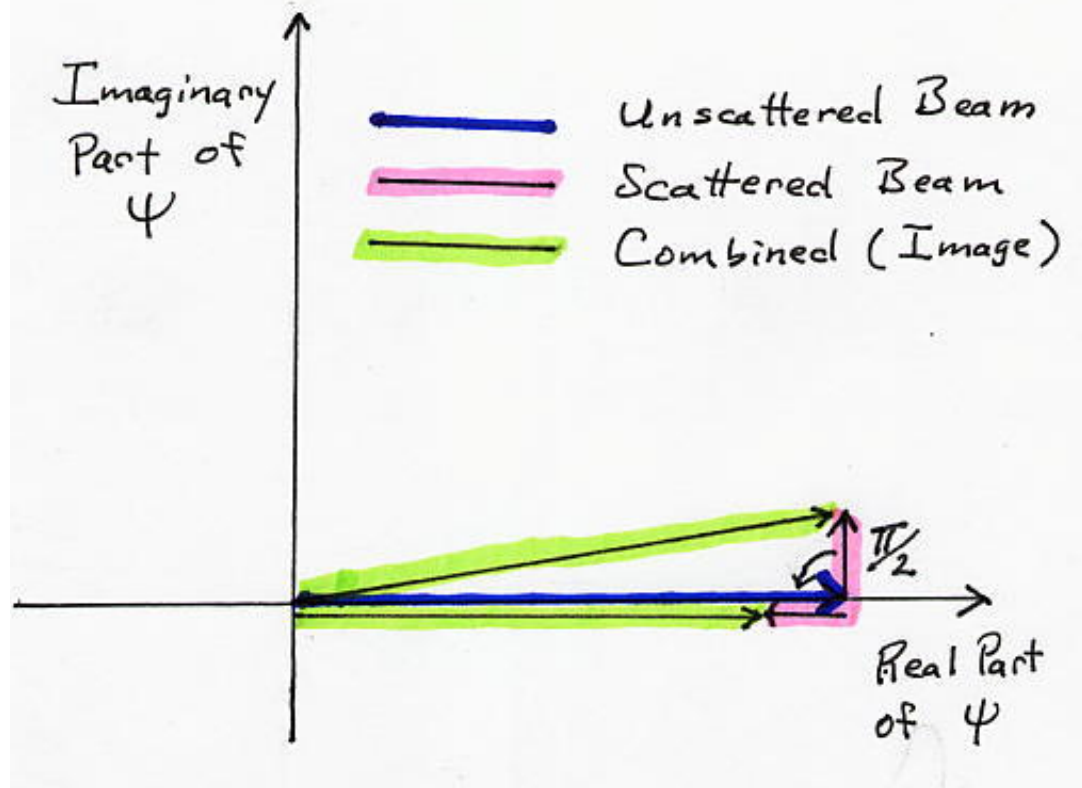
$h(\mathbf{x})$  is the *point spread function* of the image wave function – It is the inverse Fourier transform of  $H(\mathbf{s})$

# IMAGE CONTRAST REFLECTS CHANGES IN BOTH THE PHASE AND THE AMPLITUDE OF THE ELECTRON WAVES



- A SPECIMEN IS A PURE **PHASE OBJECT** IF THE **TRANSMITTED AMPLITUDE IS CONSTANT** BUT PHASE IS NOT
- A SPECIMEN IS A PURE **AMPLITUDE OBJECT** IF THE **TRANSMITTED PHASE IS CONSTANT** BUT AMPLITUDE IS NOT
- **REAL OBJECTS ARE ALWAYS MIXED**, BUT AMPLITUDE CONTRAST IS VERY WEAK IN CRYO-EM SPECIMENS

# PHASE-CONTRAST OBJECTS REQUIRE A $\pi/2$ PHASE SHIFT TO BE SEEN

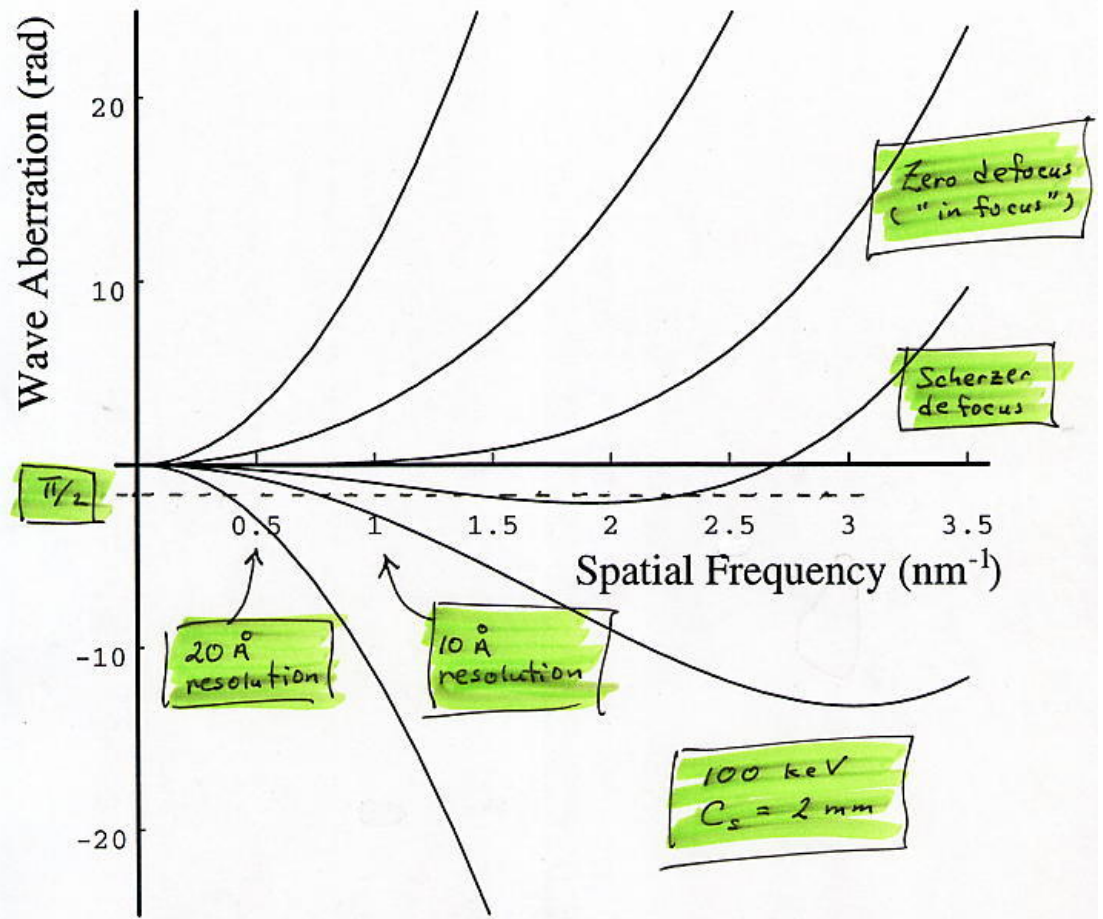


- THE SCATTERED BEAM GIVES NO CONTRAST FOR A PHASE OBJECT BECAUSE IT IS  $\pi/2$  OUT OF PHASE
- APPLYING AN ADDITIONAL  $\pi/2$  PHASE SHIFT CAN THUS PRODUCE CONSIDERABLE CONTRAST



# DEFOCUS AND SPHERICAL ABBERATION CHANGE THE PHASE OF THE SCATTERED ELECTRON WAVE

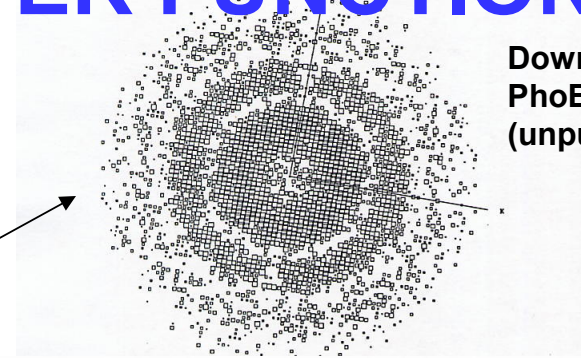
- Defocus and spherical aberration combine to change the phase
  - just as happens in the phase-contrast light microscope
- The “wave aberration” is not a uniform 90-degree phase-shift as it is in the Zernicke phase-contrast microscope, however



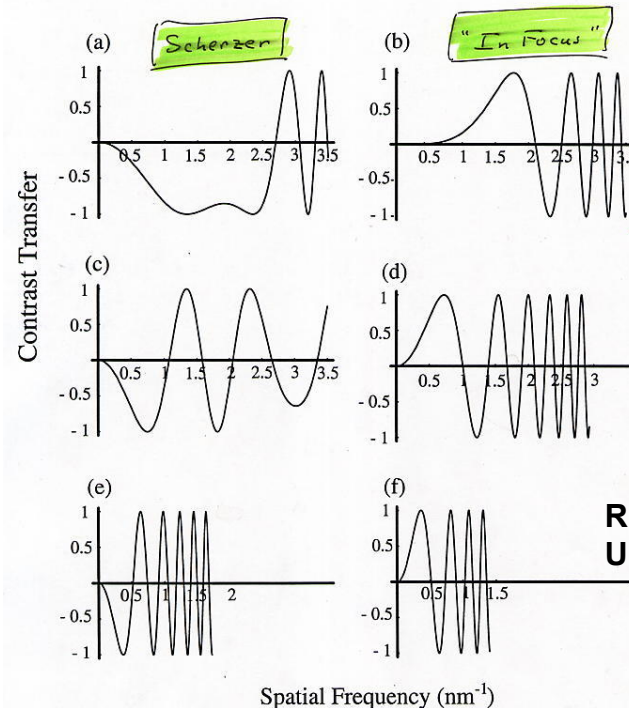
$$H(s) = \exp i\{\gamma(s)\}, \text{ and } \gamma(s) = 2\pi[C_s\lambda^3/4 s^4 - \Delta Z\lambda/2 s^2]$$

# PHASE CONTRAST IS USUALLY DESCRIBED IN TERMS OF A CONTRAST TRANSFER FUNCTION

- THE FOURIER TRANSFORM OF THE IMAGE INTENSITY IS PROPORTIONAL TO  $\text{Sin } \gamma(s) \{FT [\text{object}]\}$
- $\text{Sin } \gamma(s)$  is itself the *FT* of a point spread function for the image intensity, which is derived from  $h(x)$  mentioned in slide #7
- $\text{Sin } \gamma(s)$  IS KNOWN AS THE PHASE CONTRAST TRANSFER FUNCTION (CTF)



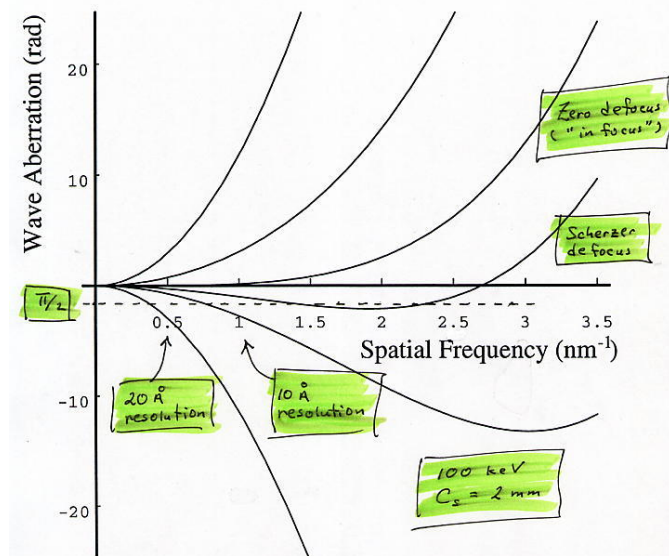
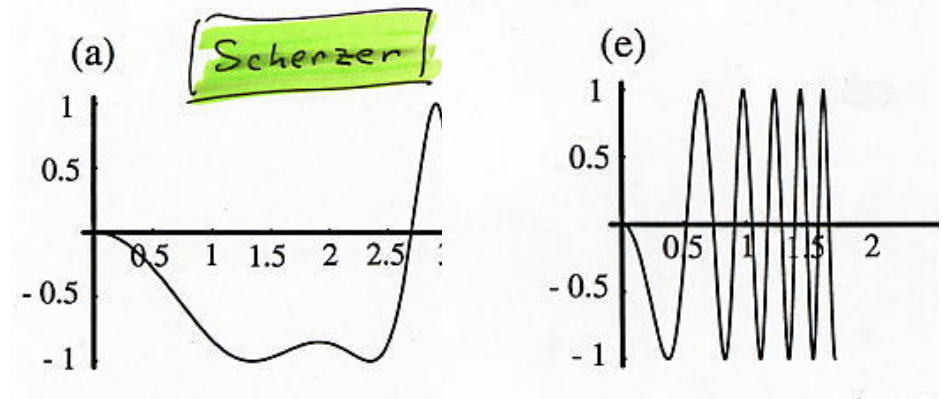
Downing & Jap  
PhoE porin image  
(unpublished)



# ONE IS TEMPTED TO USE HIGH DEFOCUS VALUES BECAUSE LOW RESOLUTION IS ALL THAT ONE CAN SEE BY EYE

- WHILE HIGH DEFOCUS MAKES IT POSSIBLE TO SEE THE OBJECT, IT ALSO CAUSES RAPID OSCILLATIONS
- THE RAPID CONTRAST REVERSALS ARE DUE TO THE STEEP INCREASE IN

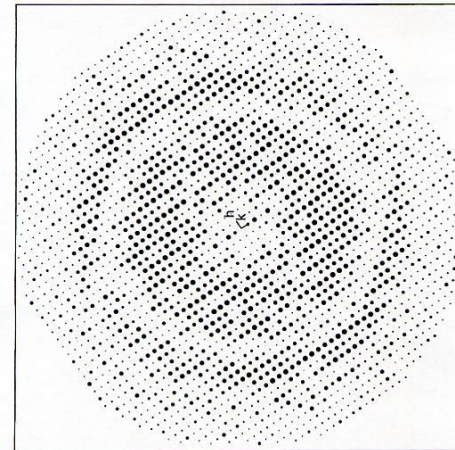
$$(\delta) \sim \pi \Delta Z \lambda s^2$$



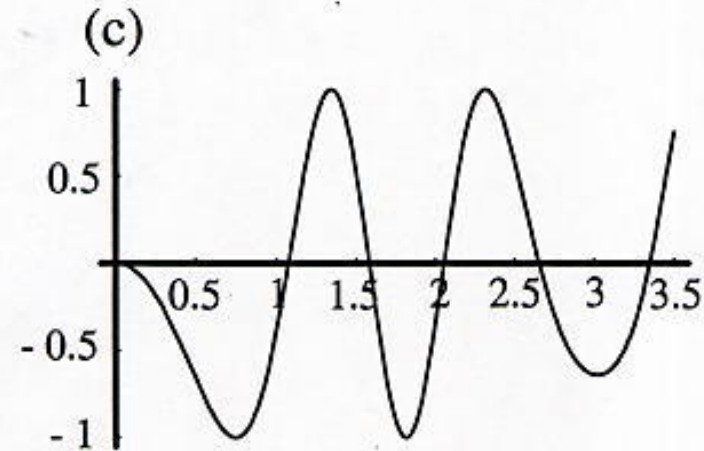
# IMAGES LOOK “ROUGHLY” LIKE A PROJECTION OF THE OBJECT

COMPUTATIONAL RESTORATION IS NECESSARY FOR QUANTITATIVE WORK, HOWEVER

- ONE MUST FIRST LOCATE THE “ZEROS” IN THE CTF
  - THEY ARE APPARENT IN THE FOURIER TRANSFORM OF THE TUBULIN CRYSTAL ON THE RIGHT
  - THEY ARE SIMILARLY APPARENT IN AREAS WITH AMORPHOUS CARBON, etc.
- SIMPLY CHANGE THE SIGN OF THE FOURIER TRANSFORM IN “EVEN” ZONES OF THE CTF
- BE AWARE THAT ASTIGMATISM INVALIDATES APPLICATION OF CIRCULAR SYMMETRY
- COMPENSATION FOR THE AMPLITUDE OF THE CTF AND THE ENVELOPE FUNCTION IS ALSO POSSIBLE DURING COMPUTATION

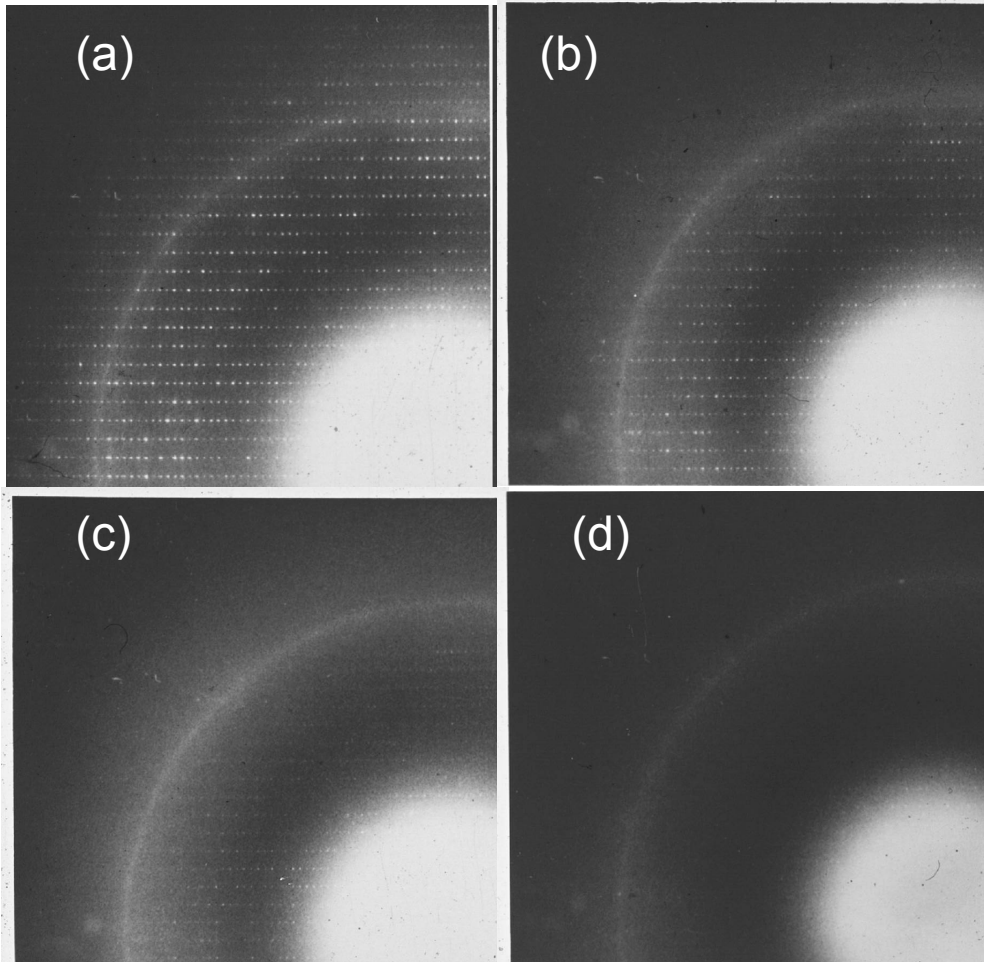


Courtesy of Ken Downing





# RADIATION DAMAGE: ELECTRONS ARE A FLUX OF IONIZING RADIATION



- **Biological macromolecules are destroyed by radiation damage**
  - Remember – there is a one-to-one connection between spots in the scattered wave and sine-functions in the image
- **Images must thus be recorded with “safe” electron exposures**
  - $< 10e/A^2$  at 100 keV
  - $< 20e/A^2$  at 300 keV
- **Bubbling sets in at doses about 3X higher than that**



# SAFE ELECTRON EXPOSURES RESULT IN INSUFFICIENT STATISTICAL DEFINITION OF HIGH-RESOLUTION FEATURES

- ALBERT ROSE DETERMINED A QUANTITATIVE RELATIONSHIP BETWEEN FEATURE SIZE AND VISUAL DETECTABILITY:

$$d C > 5 / (N)^{1/2}$$

WHERE "N" IS THE NUMBER OF QUANTA PER UNIT AREA

- FEATURES SMALLER THAN 25A MAY NOT BE DETECTABLE FOR EXPOSURES AS LOW AS 25 e/A<sup>2</sup>
- THE ONLY WAY TO OVERCOME THIS LIMITATION IS TO AVERAGE INDEPENDENT IMAGES OF IDENTICAL OBJECTS



Picture	Number of photons	High-light brightness, foot-lamberts
a	$3 \times 10^3$	$10^{-6}$
b	$1.2 \times 10^4$	$4 \times 10^{-6}$
c	$9.3 \times 10^4$	$3 \times 10^{-5}$
d	$7.6 \times 10^5$	$2.5 \times 10^{-4}$
e	$3.6 \times 10^6$	$1.2 \times 10^{-3}$
f	$2.8 \times 10^7$	$9.5 \times 10^{-3}$

# CRYSTALS MAKE IT “EASY” TO AVERAGE LARGE NUMBERS OF INDEPENDENT IMAGES

- **AVERAGING CAN BE DONE IN REAL SPACE**
- **BUT IT IS EVEN EASIER TO DO IT IN FOURIER SPACE**
  - **INFORMATION ABOUT FEATURE IN THE IMAGE THAT ARE PERIODIC MUST APPEAR IN THE DIFFRACTION SPOTS**
  - **NON-PERIODIC “NOISE” IS DISTRIBUTED UNIFORMLY AT ALL SPACIAL FREQUENCIES**
  - **YOU ELIMINATE MOST OF THE NOISE IF YOU USE JUST THE DIFFRACTION SPOTS TO DO AN INVERSE FOURIER TRANSFORM**

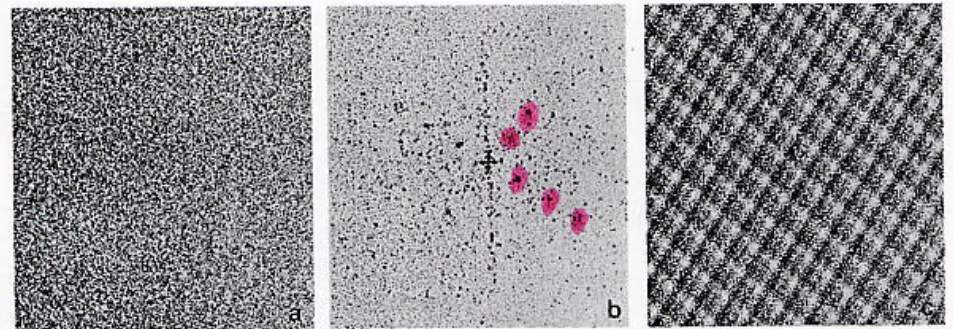


Fig. 5. (a) The Z-modulation display of a statistically noisy image of a carbon replica of an optical diffraction cross grating, recorded with an image intensifier; (b) power spectrum of the statistically noisy image; (c) the spatially averaged image.

Kuo & Glaeser (1975) Ultramicroscopy 1:53-66

- **AVERAGING A 100X100 ARRAY (i.e.  $10^4$  PARTICLES) PROVIDES THE NEEDED STATISTICAL DEFINITION REQUIRED FOR ONE VIEW (PROJECTION) AT ATOMIC RESOLUTION**

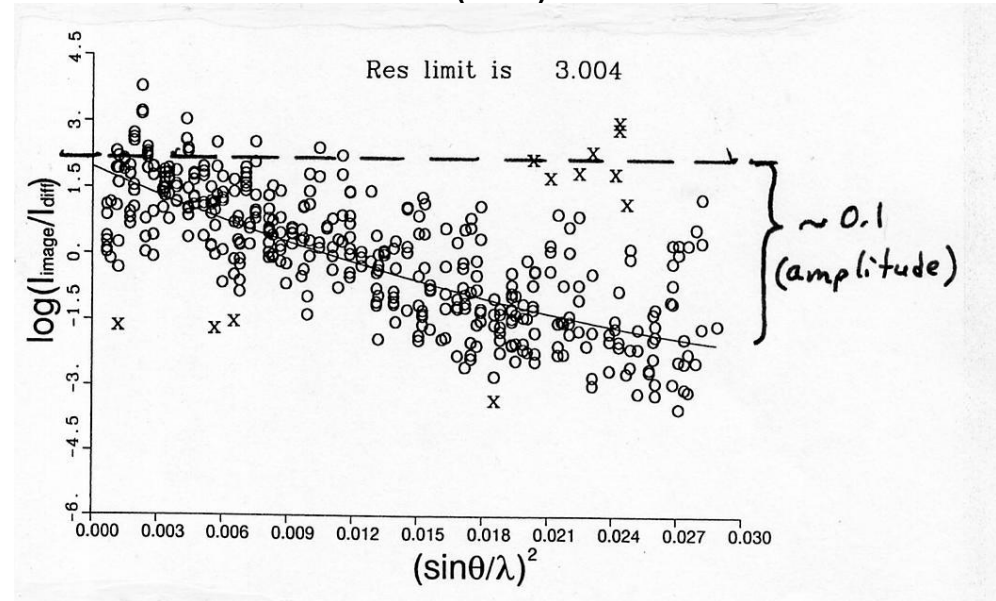
# CRYSTALS ARE NOT NECESSARY

- **ALIGN IDENTICAL PARTICLES IN IDENTICAL VIEWS BY CROSS CORRELATION**
- **CROSS CORRELATION WORKS BETTER, THE BIGGER THE PARTICLE IS**
  - BECAUSE THERE IS “MORE MASS TO BE CORRELATED”
- **PERFECT IMAGES WOULD PRODUCE ATOMIC RESOLUTION FROM ~12,000 PARTICLES AS SMALL AS  $M_r = 40,000$** 
  - INCREASE BOTH FIGURES BY 100X IF  $C = 0.1$  WHAT IT SHOULD BE [HENDERSON (1995) QUART. REV. BIOPHY.]
- **COMPUTATIONAL ALIGNMENT IS EQUIVALENT TO CRYSTALLIZATION *IN SILICO***

# MOST IMAGES CAPTURE ONLY 10% (OR LESS) OF THE SIGNAL THAT IS IN THE SCATTERED WAVE FUNCTION

- BEAM-INDUCED MOVEMENT IS THOUGHT TO BE THE CURRENT LIMITATION
- CONTRAST *CAN BE* OCCASIONALLY CLOSE TO “WHAT IT SHOULD BE” IN CURRENTLY RECORDED DATA, HOWEVER  
YONEKURA/NAMBA RESULT REQUIRED SELECTION OF PARTICLE-IMAGES THAT WERE MUCH BETTER THAN THE AVERAGE

Mitsuoka et al. (1999) J. Mol. Biol. 286:861-882



**EVEN “ROUTINE”  
CRYO-EM OF BIOLOGICAL  
MACROMOLECULES IS  
CURRENTLY *BRILLIANT***

- **Chain-trace models by 2-D electron crystallography**
- **Accurate docking of atomic models of components into large, macromolecular complexes**
- **Whole-cell tomographic imaging at ~5 nm resolution**



# THE POWER OF SINGLE-PARTICLE, REAL-SPACE AVERAGING WILL ONLY KEEP GETTING BETTER

- Automated data-collection will make it trivial to collect data sets of  $10^5$  to  $10^6$  particles
- Computer speed is keeping up with the size of data sets and the demands of higher resolution (well, at least we are trying to make it so ...)
- ***SOMEONE*** is bound to solve the problem of beam-induced movement ... (and when that happens, watch out for what cryo-EM will be able to do!)